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# Block Variable Order Step Size Method For Solving Higher Order Orbital Problems

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**Abstract.** Previous numerical methods for solving systems of higher order ordinary differential equations (ODEs) directly require calculating the integration coefficients at every step. This research provides a block multi step method for solving orbital problems with periodic solutions in the form of higher order ODEs directly. The advantage of the proposed method is, it requires calculating the integration coefficients only once at the beginning of the integration is presented. The derived formulae is then validated by running simulations with known higher order orbital equations. To provide further efficiency, a relationship between integration coefficients of various order is obtained.

## INTRODUCTION

Various real-life applications ranging from the complexity of engineering to the volatility of finance can be modeled into mathematical equations, specifically differential equations. In this research, we deal with higher order ordinary differential problems with periodic solutions.

Consider the higher order ODEs of the form

$$y^{(d)} = f(x, \tilde{Y}), \quad (1)$$

with the initial conditions  $\tilde{Y}(\alpha) = \tilde{\zeta}$ , where

$$\tilde{Y}(\alpha) = (y, y', \dots, y^{(d-1)}) \quad \text{and} \quad \tilde{\zeta} = (\zeta, \zeta', \dots, \zeta^{(d-1)}), \quad (2)$$

in the interval  $\alpha \leq x \leq \beta$ .

Referring to works by authors such as [1], [2], [3], [4] and [5], we developed a two-point block backward difference formulation with a variable order step size (VOS) code for solving higher order ODEs directly. Initially, Suleiman in [4] devised a VOS algorithm to solve stiff and nonstiff higher order ODEs directly in divided difference form known as the direct integration method (DI) which was shown to be efficient in terms of accuracy and calculation cost. The only drawback to DI is the tedious calculations of integration coefficients at every step change. Since then, the research on solving higher order ODEs directly has been expanded by authors such as [6], [7], [8], [9] and etc.

The current research, proposes a two point block method in the form of a predictor (explicit) - corrector (implicit) formulae. Because in a backward difference formulation, the integration coefficients are calculated only once at the beginning and by obtaining a relationships between the predictor and corrector of different blocks, we are able to reduce the computational costs. Next, we proceed with the formulation of the backward difference integration coefficients.

## DERIVATION OF THE PREDICT-CORRECT BACKWARD DIFFERENCE METHOD

By integrating (1) once and denoting the numbered of blocks by  $b$  we have

$$y(x_{n+1}) = y(x_n) + \int_{x_n}^{x_{n+1}} f(x, \tilde{Y}) dx, \quad b = 1, 2. \quad (3)$$

Substituting the the Newton-Gregory backward difference polynomial,

$$P_n(x) = \sum_{i=0}^{k-1} (-1)^i \binom{-s}{i} \nabla^i f_n, \quad s = \frac{x - x_n}{h} \quad (4)$$

and changing the limit of integration by replacing  $dx = hds$  gives

$$y^{(d-1)}(x_{n+1}) = y^{(d-1)}(x_n) + \int_0^b \sum_{i=0}^{k-1} (-1)^i \binom{-s}{i} \nabla^i f_n h ds \quad (5)$$

The equation above can be denoted by

$$y^{(d-1)}(x_{n+1}) = y^{(d-1)}(x_n) + h \sum_{i=0}^{k-1} \gamma_{b,1,i} \nabla^i f_n ds \quad \text{where} \quad \gamma_{b,1,i} = (-1)^i \int_0^b \binom{-s}{i} ds. \quad (6)$$

Next, we define  $G_{b,1}(t)$  as the generating function for the coefficients  $\gamma_{b,1,i}$  as follows

$$G_{b,1}(t) = \sum_{i=0}^{\infty} \gamma_{b,1,i} t^i. \quad (7)$$

By substituting  $\gamma_{b,1,i}$  from (6) into (7), the generating function  $G_{b,1}(t)$  can be written as

$$G_{b,1}(t) = \int_0^b e^{-s \log(1-t)} dt. \quad (8)$$

Solving the integration above, allows the generating function to be represented as

$$G_{b,1}(t) = - \left[ \frac{(1-t)^{-b}}{\log(1-t)} - \frac{1}{\log(1-t)} \right]. \quad (9)$$

Then, by expanding the generating function in terms of  $\gamma_{b,1,k}$  provides the following integration coefficients

$$\gamma_{b,1,k} = 1 - \sum_{i=0}^{k-1} \left( \frac{\gamma_{b,1,i}}{k-i+1} \right), \quad k = 1, 2, \dots, \quad \gamma_{b,1,0} = b. \quad (10)$$

The  $d$ th order generating function is obtained by integrating (1)  $d$  times and by mathematical induction the general solution has the following form

$$y(x_{n+1}) = y(x_n) + (bh)y'(x_n) + \dots + \frac{(bh)^{(d-1)}}{(d-1)!} y^{(d-1)}(x_n) + \frac{h^d}{d!} \sum_{i=0}^{k-1} \gamma_{b,d,i} \nabla^i f_n ds. \quad (11)$$

Similarly to  $G_{b,1}(t)$ , the  $d$ -th generating function,  $G_{b,d}(t)$  gives

$$G_{b,d}(t) = \frac{1}{(d-1)!} \left[ \frac{b^{d-1}}{\log(1-t)} - \frac{(d-1)! G_{b,d-1}(t)}{\log(1-t)} \right] \quad (12)$$



where the coefficients is independent of  $k$ . Next, is the corrector taking the form of

$$\begin{aligned}
cr y_{n+b}^{(d)} &= \sum_{i=0}^k \gamma_{b,d,i}^* \nabla_{pr}^i f_{n+b} \\
cr y_{n+b}^{(d-1)} &= y_n^{(d-1)} + h \sum_{i=0}^k \gamma_{b,d,i}^* \nabla_{pr}^i f_{n+b} \\
&\vdots \\
cr y_{n+b} &= \sum_{i=0}^{d-1} \frac{(bh)^i}{i!} y_n^{(i)} + h^d \sum_{i=0}^k \gamma_{b,d,i}^* \nabla_{pr}^i f_{n+b}.
\end{aligned} \tag{17}$$

Because of the relationship between explicit and implicit coefficients as shown in (15), the corrector can be expressed as

$$\begin{aligned}
cr y_{n+b}^{(d)} &= {}^{pr} y_{n+b}^{(d)} + \gamma_{b,0,i}^* \nabla_{pr}^i f_{n+b} \\
cr y_{n+b}^{(d-1)} &= {}^{pr} y_{n+b}^{(d-1)} + h \gamma_{b,1,i}^* \nabla_{pr}^i f_{n+b} \\
&\vdots \\
cr y_{n+b} &= {}^{pr} y_{n+b} + h^d \gamma_{b,d,i}^* \nabla_{pr}^i f_{n+b}.
\end{aligned} \tag{18}$$

The local truncation error (LTE) is represented by

$$\begin{aligned}
\widetilde{E}_k &= \gamma_{b,0,k}^* \nabla_{pr}^k f_{n+b} \\
\widetilde{E}_k^1 &= h \gamma_{b,1,k}^* \nabla_{pr}^k f_{n+b} \\
&\vdots \\
\widetilde{E}_k^d &= h^d \gamma_{b,d,k}^* \nabla_{pr}^k f_{n+b}.
\end{aligned} \tag{19}$$

by manner of Milne error estimate. As described in [14], for an efficient order and step size algorithm a suitable  $\widetilde{E}_k^{(d-p)}$  is dependent on the appropriate  $p$ . Due to the  $P_k E C_{k+1} E$  algorithm applied in this research, the asymptotic validity can be established using

$$\widetilde{E}_{k+1}^{d-p} = h^{d-p} \gamma_{b,d-1,k+1}^* \nabla^{k+1} f_{n+b}. \tag{20}$$

## NUMERICAL RESULTS

Numerical solutions for solving orbital problems in the form of second order ODEs have been investigated by authors such as [15], [16] and [17]. In the current work, the 2PBVOS method is tested with known orbital problems. Numerical results are compared against the Direct Integration (DI) method established by Suleiman in [4]. We compare the accuracy and total steps taken between both methods to determine the efficiency of the 2PBVOS method. Error calculations are done with respect to absolute error, relative error and mixed error tests (see [10]). Notations below indicates

FSTPS: failed steps,	MAXERR: the overall maximum error,
STPS: total steps,	DI: direct integration,
TOL: the tolerance used,	2PBVOS: two-point block backward difference.

### Orbit Problem 1: Two-body problem

$$\begin{aligned}
y_1''(x) &= \frac{-y_1(x)}{r}, & y_1(0) &= 1, & y_1'(0) &= 0, \\
y_2''(x) &= \frac{-y_2(x)}{r}, & y_2(0) &= 0, & y_2'(0) &= 1, \\
0 &\leq x \leq 16\pi
\end{aligned} \tag{21}$$

where

$$r = (y_1^2(x) + y_2^2(x))^{\frac{3}{2}} \quad (22)$$

with the analytical solution

$$y_1(x) = \cos x, \quad y_2(x) = \sin x. \quad (23)$$

### Orbit Problem 2: Stiefel and Bettis [18]

$$y''(x) = -y(x) + \varepsilon e^{it}, \quad y(0) = 1, \quad y'(0) = 0.9995i, \quad y \in \mathbb{C}, \quad (24)$$

$$0 \leq x \leq 1000,$$

can be written the equivalent form

$$y_1''(x) = -y_1(x) + \varepsilon \cos x, \quad y_1(0) = 1, \quad y_1'(0) = 0, \quad (25)$$

$$y_2''(x) = -y_2(x) + \varepsilon \sin x, \quad y_2(0) = 0, \quad y_2'(0) = 0.9995,$$

where  $\varepsilon = 0.001$  with the analytical solution

$$y(x) = y_1(x) + iy_2(x), \quad y_1, y_2 \in \mathbb{R}, \quad (26)$$

$$y_1(x) = \cos x + \frac{1}{2}\varepsilon x \sin x, \quad y_2(x) = \sin x + \frac{1}{2}\varepsilon x \cos x. \quad (27)$$

TOL	MTD	STPS	FSTPS	MAXERR
$10^{-2}$	DI	95	5	1.00000(0)
$10^{-2}$	2PBVOS	43	3	1.00000(0)
$10^{-4}$	DI	137	2	1.27865(-2)
$10^{-4}$	2PBVOS	81	2	9.67312(-2)
$10^{-6}$	DI	179	2	1.55778(-3)
$10^{-6}$	2PBVOS	150	2	3.62825(-3)
$10^{-8}$	DI	210	0	2.68910(-4)
$10^{-8}$	2PBVOS	329	8	7.13106(-5)
$10^{-10}$	DI	366	1	5.69993(-7)
$10^{-10}$	2PBVOS	398	15	3.67720(-6)

TABLE 1: Numerical Results for Problem 1

TOL	MTD	STPS	FSTPS	MAXERR
$10^{-2}$	DI	1176	1	1.06797(0)
$10^{-2}$	2PBVOS	922	5	1.04644(0)
$10^{-4}$	DI	2488	4	9.60862(-3)
$10^{-4}$	2PBVOS	1073	0	7.98562(-2)
$10^{-6}$	DI	3145	1	3.13723(-4)
$10^{-6}$	2PBVOS	2471	1	2.33381(-4)
$10^{-8}$	DI	3650	1	6.71619(-5)
$10^{-8}$	2PBVOS	5701	31	3.02409(-5)
$10^{-10}$	DI	6615	2	1.37306(-7)
$10^{-10}$	2PBVOS	6623	16	6.17905(-7)

TABLE 2: Numerical Results for Problem 2

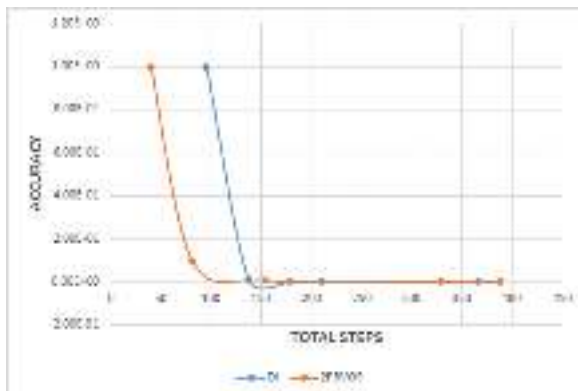


FIGURE 1: Efficiency of 2PBVOS and DI method for Problem 1

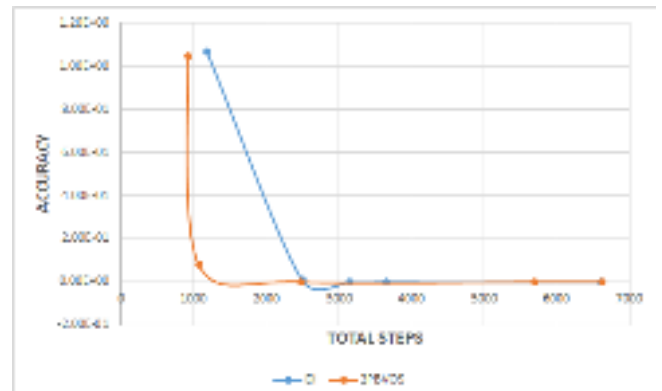


FIGURE 2: Efficiency of 2PBVOS and DI method for Problem 2

## DISCUSSION AND CONCLUSION

The division component in a divided difference formulation might effect the round off errors of the DI method especially with trigonometric solutions. This might be the cause of difference in total steps. As shown in Table 1 and 2, the 2PBVOS method requires significantly less number steps, for bigger tolerances without loss of accuracy. When dealing with finer tolerances, both methods show to be competitive.

Figure 1 and 2 represent analysis of accuracy against number of total steps. As clearly illustrated in both figures, the 2PBVOS method has the undermost curve compared to the DI method for large tolerances. Being that the undermost curve denotes the efficiency of the method (see [14]), we can conclude that the 2PBVOS method show to be more sufficient then the DI method

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